Electron Scattering

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Quasi-Elastic Scattering Coulomb Sum Rule Controversy and Future Experiment

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Studying Nuclear Structure with Electron Scattering

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Smaller, smaller

- Quest for basic building blocks of the universe
- The concept of atom by Greeks
 - atom: indivisible, fundamental building block
- Modern atoms are composed of nuclei and electrons
- Nucleus is composed of protons and neutrons
 - Oxygen atom (O) has 8 electrons and a nucleus with 8 protons and 8 neutrons
- Nucleons (protons and neutrons) are composed of quarks, gluons
 - In naïve quark model, proton has 2 u quarks and 1 d quark.

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Observing atoms, or anything else

Optical microscope

- Send light to the object
- Re-emitted light is focused/magnified by lenses
- Human eye detects the light
- Image recognized in the brain
- Resolution:
 \lambda of visible light (1000 angstrom or 100 nm)

Electron microscope

- Send electron beam to the object
- Scattered electrons are focused/magnified by coils
- Phosphor screen detects scattered electrons
- Human eye sees the image
- Resolution: \u03c0 of electron matter wave (typically 0.2 nm)

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Observing Nucleus

- Size of the nucleus: a few fm (femto meter, or fermi, 1 fm = 10⁻¹⁵m)
- Requires electron beam at higher energies (100's of MeV)

 $\hbar c = 197 \mathrm{MeV} \cdot \mathrm{fm}$

- Electron beam is prepared by accelerators
- Focusing/bending electrons are done by magnets
- Sophiscated detector system for the electrons.
- Observe number of scattered electrons as a function of angle and energy
- Reconstruction of *image* from the cross section

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Why use electrons?

One of elementary particles

- point-like without internal structure
- stable particle
- well-known properties (mass, charge, spin etc)
- well-known interaction with other elementary particles (*e.g.* quarks)
- Relatively easy to prepare (compare to other leptons or quarks)
- Easy to detect
- In general, experiments are quite *clean* compared to those with hadron beams

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Electron Scattering from a Point Charge



Rutherford-like scattering

$$\left(\frac{d\sigma}{d\Omega}\right)_{\text{point}} \equiv \left(\frac{d\sigma}{d\Omega}\right)_{\text{Mott}} = \frac{(Z\alpha)^2 \cos^2 \frac{\theta}{2}}{4E^2 \sin^4 \frac{\theta}{2}}$$

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Probing the Nucleus



Fermi's Golden Rule

$$\frac{d\sigma}{d\Omega} = \frac{2\pi}{\hbar} \left| \mathcal{M}_{fi} \right|^2 D_f$$

- *M_{fi}*: scattering amplitude
- D_f: density of the final states (or phase factor)

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Potential Scattering Amplitude

$$\mathcal{M}_{fi} = \int \psi_f^* V(\mathbf{x}) \psi_i \, d^3 \mathbf{x}$$

= $\int e^{-i\mathbf{k}_f \cdot \mathbf{x}} V(\mathbf{x}) e^{i\mathbf{k}_i \cdot \mathbf{x}} \, d^3 \mathbf{x}$
= $\int e^{i\mathbf{q} \cdot \mathbf{x}} V(\mathbf{x}) \, d^3 \mathbf{x}$

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- Plane wave approximation for incoming and outgoing electrons
- Born approximation (interact only once)

Form Factor and Charge Distribution

Using Coulomb potential from a charge distribution $\rho(\mathbf{x})$,

$$V(\mathbf{x}) = -\frac{Ze^2}{4\pi\varepsilon_0} \int \frac{\rho(\mathbf{x}')}{|\mathbf{x} - \mathbf{x}'|} d^3\mathbf{x}'$$

$$\mathcal{M}_{fi} = -\frac{Ze^2}{4\pi\varepsilon_0} \int e^{i\mathbf{q}\cdot\mathbf{x}} \int \frac{\rho(\mathbf{x}')}{|\mathbf{x} - \mathbf{x}'|} d^3\mathbf{x}' d^3\mathbf{x}$$
$$= -\frac{Ze^2}{4\pi\varepsilon_0} \int e^{i\mathbf{q}\cdot\mathbf{R}} \left[\int \frac{e^{i\mathbf{q}\cdot\mathbf{x}'}\rho(\mathbf{x}')}{|\mathbf{R}|} d^3\mathbf{x}' \right] d^3\mathbf{R}$$
$$= -\frac{Ze^2}{4\pi\varepsilon_0} \int \frac{e^{i\mathbf{q}\cdot\mathbf{R}}}{\mathbf{R}} d^3\mathbf{R} \int e^{i\mathbf{q}\cdot\mathbf{x}'}\rho(\mathbf{x}') d^3\mathbf{x}'$$
$$F(\mathbf{q}) = \int e^{i\mathbf{q}\cdot\mathbf{x}'}\rho(\mathbf{x}') d^3\mathbf{x}'$$

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Form Factor and Cross Section

For point-like particle, $\rho(\mathbf{x}') = \delta(\mathbf{x}')$ and $F(\mathbf{q}) = 1 \rightarrow$ Rutherford-*like* scattering

$$\left(\frac{d\sigma}{d\Omega}\right)_{\text{point}} \equiv \left(\frac{d\sigma}{d\Omega}\right)_{\text{Mott}} = \frac{(Z\alpha)^2 \cos^2 \frac{\theta}{2}}{4E^2 \sin^4 \frac{\theta}{2}}$$

Scattering from a charge distribution

$$\frac{d\sigma}{d\Omega} = \left(\frac{d\sigma}{d\Omega}\right)_{\mathsf{Mott}} |\boldsymbol{F}(\mathbf{q})|^2$$

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Charge Distribution in the Nucleus

- Measurement of F(q) (from cross section)
- Inverse Fourier transform gives $\rho(\mathbf{x})$



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Probing Inside the Proton

 For a target with non-zero spin - form factors for charge and magnetization

$$\begin{aligned} \frac{d\sigma}{d\Omega_{\text{lab}}} &= \left(\frac{\alpha^2 E'^2 \cos^2 \frac{\theta}{2}}{4E^3 \sin^4 \frac{\theta}{2}}\right) \left[(G_E^p)^2 + \frac{\tau}{\varepsilon} (G_M^p)^2 \right] \left(\frac{1}{1+\tau}\right) \\ \tau &\equiv \frac{Q^2}{4M^2}, \quad \frac{1}{\varepsilon} = 1 + 2(1+\tau) \tan^2 \frac{\theta}{2} \end{aligned}$$

 G_E^p distribution of charge inside the proton G_M^p distribution of magnetization inside the proton

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Electric and Magnetic Form Factor of the Proton



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New Method to Measure Proton Form Factors

- Use of *polarized* electron beam
- Measure polarization transfer to the proton

$$I_0 P_t = -2\sqrt{\tau(1+\tau)} G_E^p G_M^p \tan \frac{\theta}{2}$$

$$I_0 P_l = \frac{(E+E')}{M} \sqrt{\tau(1+\tau)} (G_M^p)^2 \tan^2 \frac{\theta}{2}$$

$$I_0 = (G_E^p)^2 + \frac{\tau}{\varepsilon} (G_M^p)^2$$

$$\frac{G_E^p}{G_M^p} = -\frac{P_t}{P_l} \frac{(E+E')}{2M} \tan \frac{\theta}{2}$$

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Surprise!



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Why?

- Same object, two different way of looking
- Current smoking gun \rightarrow Rosenbluth separation
 - Initial assumption: Born approximation (one photon exchange)
 - At large Q², two photon exchange needs to be considered



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One Possibility



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Protons Inside the Nucleus

- Quasi-elastic scattering scattering from the bound protons inside the nucleus
- Useful tool to investigate nucleon properties inside the nucleus



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Cross Section

$$\frac{d^2\sigma}{d\Omega d\omega} = \sigma_{\text{Mott}} \left[\frac{Q^4}{q^4} R_L(q,\omega) + \frac{Q^2}{2q^2} \frac{1}{\varepsilon} R_T(q,\omega) \right]$$

- ▶ $R_L(q, \omega), R_T(q, \omega)$: Response functions
- Analogy of G_E^p and G_M^p of the free proton
- *R_L(q, ω)* characterizes charge interaction in the nucleus.
- Coulomb Sum

$$S_L(q) = \int_{\omega_{\rm el}^+}^{\infty} d\omega \frac{R_L(q,\omega)}{Z\tilde{G_E}^2(Q^2)}$$

$$\tilde{G_E}^2(Q^2) = \left([G_E^p(Q^2)]^2 + (N/Z) [G_E^n(Q^2)]^2 \right) \frac{1 + Q^2 / 4M^2}{1 + Q^2 / 2M^2}$$

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Coulomb Sum Rule

- By definition, $S_L(q) = 1$ for the free proton
- Saturation of the Coulomb Sum $S_L(q) \rightarrow 1$ at sufficiently large q
- Deviation of the Coulomb Sum Rule
 - at small q Nucleon-nucleon long-range correlations and Pauli blocking
 - at large q Short range correlations and modification of the free nucleon electromagnetic properties inside the nuclear medium
- Nuclear density dependence (⁴He to ²⁰⁸Pb)
- Related to chiral symmetry restoration in dense nuclear medium

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Measurements

- For the past twenty years, a large experimental program at Bates, Saclay and SLAC
- Limited kinematic coverage in q and ω due to machine limitations





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Controversy



- Early data shows significant quenching of the CSR.
- With the addition of forward angle data, Bates claims no significant quenching.
- Saclay new analysis claims that quenching persists.

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Controversy and Future Experiment



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New Proposed Experiment at JLab



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- Electron Scattering Experiments
 - One of the main tools in nuclear physics research
 - Charge density of the nucleus Comparison with theoretical model
 - Structure of the proton (charge and magnetization)
- New challenge
 - Surprising results on G_E^p/G_M^p from JLab
 - Two photon exchange effect more study in theory and experiment
- Nucleons inside the nucleus
 - Change or no change, that is the question.
 - Controversial results on the Coulomb Sum
 - Proposed experiment will give a definitive answer
- Probing quarks
 - Existence and interaction of quarks
 - Study of the quark spins

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